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Research Article

Forward Kinematic of a Sphere Considering Slipping and Motion Analysis in Three Rollers

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ABSTRACT

Compared to mobile robots equipped with multiple rollers, spherical robots can move in all directions and are superior in durability and in their ability to climb steps. Slippage between balls and rollers is a significant problem in friction drives. However, previously established rollerdriven ball kinematics model considers sliding on only two constraining rollers. In this research, we developed it, proposed a motion model of sphere with three-constraint rollers, and developed a mathematical model that simulates the angular velocity vector of the sphere and the slip vector at each contact point. And we considered the existence of an angular velocity vector of sphere adapted three constraint rollers from the viewpoint of forward kinematics and succeed demonstration of the trajectory of the endpoint of the angular velocity vector and slip velocity vector of sphere.

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1. Introduction

Spherical robots are widely used in robotics, such as multi-finger fingertip mechanism for hand robots, actuator transmission mechanisms for omnidirectional movement, and driving mobile robots. They are also used as driving rollers for omnidirectional movement mechanisms; there are various arrangements and spherical structures depending on the requirements of the movement mechanism. Figure 1 shows the roller contact type for the number of actuators (N_R) per sphere.

In the case of $N_R = 2$, ACROBAT-S [1] (Figure 1(a)) and wheel chair [2] (Figure 1(b)) with sphere kinematics are developed for transportation service. The omnidirectional locomotion condition is that two rotational axes contact point with sphere are arranged along the great circle. Thus, the angular velocity vector of the sphere lies on a common plane parallel to the great circle of sphere. In this way, the roller arrangement conditions for omni directional spherical mobile robot are derived [3]. Furthermore, in this situation, the angular velocity vector of the sphere has two degrees of freedom. Using this theory, spherical robot transfer problem is considered in [4]-[7], as models don't consider slipping.

As shown in Figure 1(c), the examples of the ballholding mechanism are Musashi150 [8], RV-infinity [9], and NuBot [10], which are designed to transport ball. Most teams in the RoboCupMSL(Middle-Size-League) adapts a ball-holding mechanism with two rollers on the upper hemisphere to control the rotational motion of the ball. In most developments, roller arrangement that causes slippage (when the roller rotational axis don't have slip in along a great circle) is adopted. Now there is slippage at contact point between the roller and the ball. It utilizes the ability to hold the ball due to frictional force. However, in the absence of a suitable mathematical model, the roller arrangement is heuristically determined experimentally.

In a previous study, we used two constraint rollers that allowed the slipping to derive a mathematical

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model for spherical rotational motion for forward kinematics [11]. This model is included in the kinematics of [3]. Furthermore, we employed experiment [12] to validate the model of [11].

In the case of $N_R = 3$, omnidirectional-wheeled mobile platform (OWMP) [13] has three-constraint rollers (Figure 1(d)), while a ball-balanced robot [14] has three unconstraint rollers (Figure 1(e)). In the example of an adapted unconstraint three rollers, ball-valanced robot can move in omni-direction [15]. Furthermore, it attempts to carry the load using couple of it ([15] and [16]). Ball valanced robot (Figure 1(e)) [14] has an adapted roller in the upper hemisphere, whereas Atras sphere [17] has an adapted roller in the lower hemisphere and can control sphere.

Although there are three rollers, the sphere rotational dimensions are different because of the roller structure (constraint type or unconstraint type). Each constrained roller has less rotational diversity than the unconstrained rollers. However, the holding force is stronger than that of the unconstrained roller. The stability of the sphere rotational motion is higher in the case of three rollers than in the case of two rollers. OWMP [13] is used kinematic with roller arrangement restricted to the equator; we extended this kinematics to an arbitrary roller's arrangement discussion.

In this study, we modified the previously developed kinematic model [11] for the case of three-constraint rollers. And presented a novel mathematical model



Figure 1. A pattern of roller arrangement for sphere mobile robot driven by multiple rollers.

about sphere rotational motion using three-constraint rollers that allows for slipping. Additionally, we simulate the angular velocity vector of the sphere, the sphere mobile speed vector, slip vector between the roller and the ball and the slip speed using forward kinematics. The rest of this study is as follows: Chapter 2 discusses the existing space of angular velocity vectors on rollers and drives the sphere's kinematics with three rollers. Chapter 3 conducted the simulation. Finally, we present the summary and future works.

2. The sphere forward kinematics for three constraint rollers

In this chapter, we introduce the angular velocity vector of the sphere to model the sphere's rotational motion geometrically.

2.1 Consideration of existence of angular velocity vector of the sphere

(A) Case of single-constraint driving roller

Table 1 shows variables list for this study. Figure 2(a) shows existence space of angler velocity vector of sphere in a single-roller. The center $\mathbf{0}$ of a sphere with radius r is fixed as the origin of the coordinate system $\Sigma - xyz$. The constraint roller i have the center of gravity \mathbf{R}_i and \mathbf{q}_i denotes contact point between roller i and sphere. η_i denotes the unit vector along the rotational axis of the constraint roller. $\mathbf{0}$, \mathbf{q}_i and \mathbf{R}_i are arranged on common line. $\boldsymbol{\omega}$ denotes the angular velocity vector of the sphere. v_i denotes the peripheral speed of the constraint roller i at \mathbf{q}_i on sphere. $\mathcal{V}_i^S(=\boldsymbol{\omega} \times \mathbf{q}_i)$ denote the velocity vector of the sphere with respect to \mathbf{q}_i .



Figure 2 (a) The existence of sphere angler velocity vector in the case of a single-constraint roller. (b) Roller's axis vector η_i at contact point q_i on the sphere

Table 1 Variables list

$\Sigma - xyz$	Three-dimensional coordinate system fixed the center of sphere
$\langle x, y \rangle$	Inner product with respect to \boldsymbol{x} and \boldsymbol{y}
$\ x\ $	Norm of vector \boldsymbol{x}
span{ x , y }	Plane spanned by \boldsymbol{x} and \boldsymbol{y}
0	Origin of $\Sigma - xyz$ (Center of sphere)
q_i	Contact point from sphere and roller
$\boldsymbol{\eta}_i$	Unit vector of constraint-roller along the
ω	rotational axis direction Sphere angular velocity vector
R_i	Center of gravity of constraint roller
v_i^R	Velocity vector of constraint roller i at q_i
v_i^S	Velocity vector of sphere at q_i
$\zeta_{i,k}$	Slip velocity of the sphere with respect to v_i^R in k_i th pattern of roller arrangement $(k = 1, 2, 3, 4)$
\boldsymbol{V}_k	Sphere mobile speed in <i>k</i> -th pattern of roller arrangement ($k = 1,2,3,4$)
$\boldsymbol{\omega}_k$	Sphere angular velocity vector in <i>k</i> -th pattern of roller arrangement ($k = 1,2,3,4$)
\boldsymbol{e}_i	Unit normal vector along v_i^R
V	Sphere robot velocity vector on xy-plane
ψ_i	Roller arrangement angle that rotates
	counterclockwise with respect to
$l_i(v_i)$	Existence set of end point of $\boldsymbol{\omega}$ with respect to
12:	v_i Perinheral speed of constraint roller <i>i</i>
r	Sphere robot radius
,	
φ	Sphere robot direction
d_i	the minimal distances between point and line



Figure 2 (a) The existence of sphere angler velocity vector in the case of a single-constraint roller. (b) Roller's axis vector η_i at contact point q_i on the sphere



Figure 3 Slip adjudication by the location $l_1(v_1)$, $l_2(v_2)$ and $l_3(v_3)$. (a). A pair $\{v_1, v_2, v_3\}$ exists such that $l_1(v_1)$, $l_2(v_2)$ and $l_3(v_3)$ have points in common. (b). A pair $\{v_1, v_2, v_3\}$ exists such that $l_1(v_1)$, $l_2(v_2)$ and $l_3(v_3)$ have no points in common.

 v_i^R denotes velocity vector of the roller. e_i denotes unit normal vector along v_i^R . Because of $e_i \in$ span $\{q_i, \eta_i\}$, e_i and v_i^R are satisfy $v_i = \langle v_i^R, e_i \rangle$. Thus. Nonslip condition is $v_i^S = v_i^R$ and ω can be satisfied as follow.

$$\boldsymbol{v}_i^{\boldsymbol{S}} = \boldsymbol{\omega} \times \boldsymbol{q}_i \tag{1}$$

 $e_i \in \operatorname{span}\{q_i, \eta_i\}$ is the unit normal vector along v_i^R . Using $v_i = \langle v_i^R, e_i \rangle$ ($v_i^S = v_i^R$: nonslip condition) and Eq. (1), v_i is represented as follows:

$$v_i = \langle v_i^R, e_i \rangle = -r \langle \eta_i, \omega \rangle$$
 (2)

Thus, $\boldsymbol{\omega}$ can be satisfied as Eq. (3).

$$\langle \boldsymbol{\eta}_i, \boldsymbol{\omega} \rangle = -\frac{v_i}{r}$$
 (3)

 $\boldsymbol{\omega}$ must be on **span**{ $\boldsymbol{\eta}_i, \boldsymbol{q}_i$ } and can be represented as a following line set $l_i(v_i)$ that is parallel to \boldsymbol{q}_i and passes through the end point of $-(v_i/r)\boldsymbol{\eta}_i$.

$$l_i(v_i) = \left\{ \boldsymbol{\omega} \middle| \left(-\frac{v_i}{r} \right) \boldsymbol{\eta}_i + t(1/r) \boldsymbol{q}_i, t \in \mathbb{R} \right\}$$
(4)

Furthermore, we set the roller contact point and rotational axis on the sphere. Contact point q_i is defined as a polar coordinate as follows:

$$\boldsymbol{q}_{i} = \begin{bmatrix} \cos \theta_{2,i} \cos \theta_{1,i}, \cos \theta_{2,i} \sin \theta_{1,i}, \sin \theta_{2,i} \end{bmatrix}^{T}$$
(5)

As shown in Figure 2(b), η_i has a starting point at R_i (O, q_i and R_i are on the same line). R_i is located on the plane π_i parallel to the tangent plane of the sphere at q_i . We put pair of vectors as normal orthogonal bases { X_i , Y_i } on π_i at R_i .

$$\boldsymbol{X}_{i} = \begin{bmatrix} -\sin\theta_{1,i} \\ \cos\theta_{1,i} \\ 0 \end{bmatrix}, \quad \boldsymbol{Y}_{i} = \begin{bmatrix} -\sin\theta_{2,i}\cos\theta_{1,i} \\ -\sin\theta_{2,i}\sin\theta_{1,i} \\ \cos\theta_{2,i} \end{bmatrix}$$
(6)

Thus, η_i is linear combination of Eq. (6) and rotates counterclockwise with respect to ψ_i .

$$\boldsymbol{\eta}_i = \boldsymbol{X}_i \cos \psi_i + \boldsymbol{Y}_i \sin \psi_i \tag{7}$$

(B) Case of three-constraint driving rollers

As shown in Figure 3, the locations between $l_1(v_1)$, $l_2(v_2)$, and $l_3(v_3)$, which depend on parameters v_1 , v_2 and v_3 , are used to determine the rotational axis of the sphere.

If a pair of v_1 , v_2 , and v_3 exists such that $l_1(v_1)$, $l_2(v_2)$, and $l_3(v_3)$ have common points, the endpoint of $\boldsymbol{\omega}$ can be uniquely determined using Eq. (4) (as i = 1,2,3) (Figure 3(a)). Based on Eq. (4) (as i = 1,2,3), $\boldsymbol{\omega}$ must be on $\operatorname{span}\{q_1, \eta_1\} \cap \operatorname{span}\{q_2, \eta_2\} \cap$ $\operatorname{span}\{q_3, \eta_3\}$. However, if a pair of v_1, v_2 , and v_3 exists such that $l_1(v_1), l_2(v_2)$, and $l_3(v_3)$ have common no points, slip can occur (Figure 3(b)). The sphere rotational axis is defined with respect to the arbitrary parameters v_1 , v_2 , and v_3 .

2.2. Calculation method of optimal point in sum of the squared distances

In this section, we calculate the optimal point $Q_0 = (x_0, y_0, z_0) (\in \mathbb{R}^3)$, which is determined such that the sum of the squared distances between $Q = (x, y, z) (\in \mathbb{R}^3)$ and $l_i(v_i)(i = 1, 2, 3)$ is minimized.

As shown in Figure 3(b), d_i denote the distances between Q and $l_i(v_i)$ in each line $l_i(v_i)$. it is represented as follows (See Appendix (A)):

$$d_{i} = \left\| \left(-\frac{v_{i}}{r} \right) \boldsymbol{\eta}_{i} + \frac{\langle \boldsymbol{P}_{i}, \boldsymbol{Q} \rangle}{r^{2}} \boldsymbol{P}_{i} - \boldsymbol{Q} \right\|$$
(8)

Therefore. Sum of the squared distances is represented as follow:

$$L(x, y, z) = d_1^{2} + d_2^{2} + d_3^{2}$$
(9)

 $(x, y, z) = (x_0, y_0, z_0)$ such that L(x, y, z) is minimal value $A_{10} - B_8^2/4B_2$ is satisfy as following value (x_0, y_0, z_0) (See Appendix(B)).

$$z_0 = -\frac{C_9}{2C_3}$$
(10)

$$y_0 = \frac{B_5 C_9 - 2C_3 A_8}{4B_2 C_3} \tag{11}$$

$$x_{0} = \frac{1}{8A_{1}B_{2}C_{3}} (-A_{4}B_{5}C_{9} + 2A_{4}B_{8}C_{3} + 2A_{6}B_{2}C_{9} - 4A_{7}B_{2}C_{3})$$
(12)

Where

$$C_3 = B_3 - \frac{B_5^2}{4B_2}, C_9 = B_9 - \frac{B_5 B_8}{2B_2}, C_{10} = B_{10} - \frac{B_8^2}{4B_2}$$
(13)

Where

$$B_{2} = A_{2} - \frac{A_{4}^{2}}{4A_{1}}, B_{3} = A_{3} - \frac{A_{6}^{2}}{4A_{1}}, B_{5} = A_{5} - \frac{A_{4}A_{6}}{2A_{1}}$$
$$B_{8} = A_{8} - \frac{A_{4}A_{7}}{2A_{1}}, B_{9} = A_{9} - \frac{A_{6}A_{7}}{2A_{1}}, B_{10} = A_{10} - \frac{A_{7}^{2}}{4A_{1}}$$
(14)

Where

$$A_{10} = L(0,0,0)$$

$$A_{1} = \frac{1}{2} (L(-1,0,0) + L(1,0,0) - 2L(0,0,0))$$

$$A_{2} = \frac{1}{2} (L(0,-1,0) + L(0,1,0) - 2L(0,0,0))$$

$$A_{3} = \frac{1}{2} (L(0,0,-1) + L(0,0,1) - 2L(0,0,0))$$

$$A_{7} = \frac{1}{2} (L(1,0,0) - L(-1,0,0))$$

$$A_{8} = \frac{1}{2} (L(0,1,0) - L(0,-1,0))$$

$$A_{9} = \frac{1}{2} (L(0,0,1) - L(0,0,-1))$$

$$A_{4} = \frac{1}{4} (L(1,1,0) - L(1,-1,0) - L(-1,1,0) + L(-1,-1,0))$$

$$A_{5} = \frac{1}{4} (L(0,1,1) - L(0,1,-1) - L(0,-1,1) + L(0,-1,-1))$$

$$A_{6} = \frac{1}{4} (L(1,0,1) - L(-1,0,1) - L(1,0,-1) + L(-1,0,-1))$$

2.3 Forward kinematics of sphere rotational motion

We defined forward kinematics as input $v_i(i = 1, 2, 3) \rightarrow$ output V. The measured sphere robot direction angle from x-axis denotes φ and defined interval as $0^\circ \le \varphi < 360^\circ$. The sphere mobile velocity vector on the xy-plane denotes V. Because that $\boldsymbol{\omega} = [\omega_x, \omega_y, \omega_z]^T$ is perpendicular to V, The forward kinematics is represented as follows:

$$\|\boldsymbol{V}\| = r\sqrt{\omega_x^2 + \omega_y^2} \tag{16}$$

φ

$$= \begin{cases} \cos^{-1} \left[\omega_y / \sqrt{\omega_x^2 + \omega_y^2} \right] (\omega_x < 0) \\ 360^\circ - \cos^{-1} \left[\omega_y / \sqrt{\omega_x^2 + \omega_y^2} \right] (\omega_x \ge 0) \end{cases}$$
(17)

3. Simulation

This chapter presents the simulation results, including the trajectory of the endpoint of the sphere angler velocity vector $\boldsymbol{\omega}_k$, the sphere mobile velocity vector \boldsymbol{V}_k , and the slip vector $\boldsymbol{\zeta}_{i,k}$ and slip speed $\|\boldsymbol{\zeta}_k\|$ in *k*-th (k =1,2,3,4) roller arrangement patterns in the case in which a regular triangle ($\theta_{1,1}, \theta_{1,2}, \theta_{1,3}$) = (30°, 150°, 270°) and $\psi_i = 0^\circ$ (i = 1,2,3) are fixed and $\theta_{2,i}$ is parameter as $0^\circ \sim 30^\circ$. The patterns are set up by $\theta_{2,i}$ (i = 1,2,3,4) as follows: Pattern I ($k = 1, \theta_{2,i} = 0^\circ$), Pattern II ($k = 2, \theta_{2,i} = 10^\circ$), Pattern III ($k = 3, \theta_{2,i} =$ 20°), Pattern IV ($k = 4, \theta_{2,i} = 30^\circ$).

As input roller speed, we define function $v_1(\varphi) = \sin(\varphi + 240^\circ)$, $v_2(\varphi) = \sin(\varphi + 120^\circ)$ and $v_3(\varphi) = \sin \varphi$ such that output: $||V_1|| = 1$ [m/s].

As output, ω_k , V_k , $\zeta_{1,k}$, $\zeta_{2,k}$, and $\zeta_{3,k}$ (k = 1,2,3,4) were indicated, such as Pattern I [k = 1; green curve], Pattern II [k = 2; red curve], Pattern III [k = 3; blue curve], and Pattern IV [k = 4; violet curve] (Figure 4-9). They were calculated using Eqs. (10)-(12) and (16), respectively.

As shown in Figure 4, ω_k (k = 1,2,3,4) draws circle trajectories and gets a small radius in turn.

As shown in Figure 5, V_k (k = 1,2,3,4) draws circle trajectories and gets a common center and small radius in turn.

As shown in Figure 6-8, $\zeta_{i,1}$ is the only origin ($\zeta_{i,1} = 0$) due to Pattern I (nonslip case), but $\zeta_{i,2}$, $\zeta_{i,3}$, and $\zeta_{i,4}$ draw ellipsoid trajectories and get a large radius in turn. In this way, V_k (or ω_k) and $\zeta_{i,k}$ are trade-off effects.

As shown in Figure 9, because Pattern I [k = 1: nonslip case], $\|\zeta_{1,1}\|$, $\|\zeta_{2,1}\|$, and $\|\zeta_{3,1}\| = 0$ [m/s].

Three functions $\|\boldsymbol{\zeta}_{1,k}\|$, $\|\boldsymbol{\zeta}_{2,k}\|$, and $\|\boldsymbol{\zeta}_{3,k}\|$ with respect to φ are cyclical functions of $2\pi/3$.

 $\|\boldsymbol{\zeta}_{1,k}\|$ (k = 2, 3, 4) have minimal values of 0.03, 0.10, and 0.19 [m/s], respectively, when $\varphi = \pi/6, 7\pi/2$

6[rad]. $\|\boldsymbol{\zeta}_{1,k}\|(k=2,3,4)$ have maximal values of 0.16, 0.30, and 0.40 [m/s], respectively, when $\varphi = 2\pi/3, 5\pi/3$.

 $\|\zeta_{2,k}\|(k=2,3,4)$ has minimal values of 0.03, 0.10, and 0.19 [m/s], respectively, when $\varphi = 5\pi/6, 11\pi/6$.

 $\|\zeta_{2,k}\|(k=2,3,4)$ has maximal values of 0.16, 0.30, and 0.40 [m/s], respectively, when $\varphi = \pi/3, 4\pi/3$.

 $\|\zeta_{3,k}\|(k = 2, 3, 4)$ has minimal values of 0.03, 0.10, and 0.19 [m/s], respectively, when $\varphi = \pi/2, 3\pi/2$, and $\|\zeta_{3,k}\|(k = 2, 3, 4)$ has maximal values of 0.16, 0.30, and 0.40 [m/s], respectively, when $\varphi = 0, \pi, 2\pi$

In this way, the direction case in which slip speed is maximal is perpendicular to the direction case in which slip speed is minimal.

4. Conclusion

In this study, we considered the existence of an angular velocity vector of sphere driven by three rollers from the viewpoint of forward kinematics and demonstrated the trajectory of the endpoint of the angular velocity vector and slip velocity vector.

In the future research, we want to conduct the inverse kinematics and consider existence space for input rollers speed.



Figure 4 The trajectory of end point of angular velocity vector $\boldsymbol{\omega}_k(k = 1,2,3,4)$.

0.2

Figure 7 The trajectory of slip vector

x-component of ξ_2 [m/s]

from sphere and roller $\zeta_{2,k}(k =$

0.4 -0.4



0.2

Figure 8 The trajectory of slip vector

nent of $\xi_3 \, [m/s]$

from sphere and roller $\zeta_{3,k}(k =$

0.4 -0.4

component of ξ_3 [m/s]

ent of Es [m.

0.1

-0.1

-0.2

1,2,3,4).

-0.2



Figure 6 The trajectory of slip vector from sphere and roller $\zeta_{1,k}(k = 1,2,3,4)$.



Figure 9 The slip speed of slip vector from sphere and roller $\|\boldsymbol{\zeta}_{1,k}\|, \|\boldsymbol{\zeta}_{2,k}\|, \|\boldsymbol{\zeta}_{3,k}\|$ (k = 1,2,3,4).

Therefore.



1,2,3,4).

z-component of ξ_2 [m/s

0.1

0

-0.1

-0.2

-0.2

(A) Calculation of distance between Q and $l_i(v_i)$

As shown in Figure 10, It is determined distance d_i such that $l_i(v_i) \perp \overrightarrow{QQ_i}$ ($Q_i \in l_i(v_i)$, $Q \in \mathbb{R}^3$). Using Eq. (4) and inner product expand formula,

0.2

mponent of ξ_2 [m/s]

$$0 = \langle \boldsymbol{q}_{i}, \boldsymbol{Q}_{i} - \boldsymbol{Q} \rangle$$

= $\langle \boldsymbol{q}_{i}, \left(-\frac{v_{i}}{r}\right) \boldsymbol{\eta}_{i} + \frac{t}{r} \boldsymbol{q}_{i} - \boldsymbol{Q} \rangle$ (A,1)
= $\left(-\frac{v_{i}}{r}\right) \langle \boldsymbol{q}_{i}, \boldsymbol{\eta}_{i} \rangle + \frac{t}{r} \langle \boldsymbol{q}_{i}, \boldsymbol{q}_{i} \rangle - \langle \boldsymbol{q}_{i}, \boldsymbol{Q} \rangle$ (A,2)

From $\langle \boldsymbol{q}_i, \boldsymbol{\eta}_i \rangle = 0$ and $\langle \boldsymbol{q}_i, \boldsymbol{q}_i \rangle = r^2$, Eq. (A,2) can be simplified, and we obtain the value *t* that minimizes the distance between \boldsymbol{Q} and $l_i(v_i)$ as follows:

$$\Leftrightarrow 0 = rt - \langle \boldsymbol{q}_i, \boldsymbol{Q} \rangle \Leftrightarrow t = \langle \boldsymbol{q}_i, \boldsymbol{Q} \rangle / r$$
 (A,3)

Substituting $\overrightarrow{QH_{i}} = \overrightarrow{OH_{i}} - \overrightarrow{OQ}$ with Eq. (A,3),

$$\overrightarrow{\boldsymbol{QH}_{i}} = \left(-\frac{v_{i}}{r}\right)\boldsymbol{\eta}_{i} + \frac{\langle \boldsymbol{q}_{i}, \boldsymbol{Q} \rangle}{r^{2}}\boldsymbol{q}_{i} - \boldsymbol{Q}$$
(A,4)





First, based on the coefficients $A_1 \sim A_{10}$, $L_1(x, y, z)$ is represented as a quadratic equation with respect to, y and z as follows:

$$L_1(x, y, z) = A_1 x^2 + A_2 y^2 + A_3 z^2$$
(B,1)
+A_4 xy + A_5 yz + A_6 zx + A_7 x + A_8 y + A_9 z + A_{10}

Based on Eq. (B,1), the following homogeneous equation is conducted. By solving Eq. (B,2), We obtain A_1 , A_2 , A_3 , A_7 , A_8 , A_9 and A_{10} .

$$L(0,0,0) = A_{10}$$

$$L(1,0,0) = A_1 + A_7 + A_{10}$$

$$L(-1,0,0) = A_1 - A_7 + A_{10}$$

$$L(0,1,0) = A_2 - A_8 + A_{10}$$

$$L(0,-1,0) = A_2 - A_8 + A_{10}$$

$$L(0,0,1) = A_3 - A_9 + A_{10}$$

$$L(0,0,-1) = A_3 - A_9 + A_{10}$$

Using Eq. (B,1), following homogeneous equation is conducted. By solving Eq. (B,3), we get A_4 , A_5 and A_6 .

$$L(1,1,0) = A_1 + A_2 + A_4 + A_7 + A_8 + A_{10}$$

$$A_4 = L(1,1,0) - (A_1 + A_2 + A_7 + A_8 + A_{10})$$

$$A_5 = L(0,1,1) - (A_2 + A_3 + A_8 + A_9 + A_{10})$$

$$A_6 = L(1,0,1) - (A_1 + A_3 + A_7 + A_9 + A_{10})$$

(B,3)

At first, Eq. (B,1) is represented as follows by completing the square with respect to x.

$$L_{2}(x, y, z) = A_{1}\left(x + \frac{A_{4}y + A_{6}z + A_{7}}{2A_{1}}\right)^{2}$$
(B,4)
+ $B_{2}y^{2} + B_{3}z^{2} + B_{5}yz + B_{8}y + B_{9}z + B_{10}$

At second, Eq. (B,4) is represented as follows by

completing the square with respect to y and z.

$$L_{3}(x, y, z) = A_{1}\left(x + \frac{A_{4}y + A_{6}z + A_{7}}{2A_{1}}\right)^{2}$$
(B,5)
+ $B_{2}\left(y + \frac{B_{5}z + B_{8}}{2B_{2}}\right)^{2} + C_{3}\left(z + \frac{C_{9}}{2C_{3}}\right)^{2} + C_{10} - \frac{C_{9}^{2}}{4C_{3}}$

 $(x, y, z) = (x_0, y_0, z_0)$ such that $L_3(x, y, z)$ is minimal value $C_{10} - B_8^2/4B_2$ is satisfied as this linear equation with respect to x_0 , y_0 and z_0 .

$$\begin{cases} x_0 + \frac{A_4 y_0 + A_6 z_0 + A_7}{2A_1} = 0 \\ y_0 + \frac{B_5 z_0 + B_8}{2B_2} = 0 \\ z_0 + \frac{C_9}{2C_3} = 0 \end{cases}$$
(B,6)

Therefore. By solving Eq. (B,6), we get it.

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